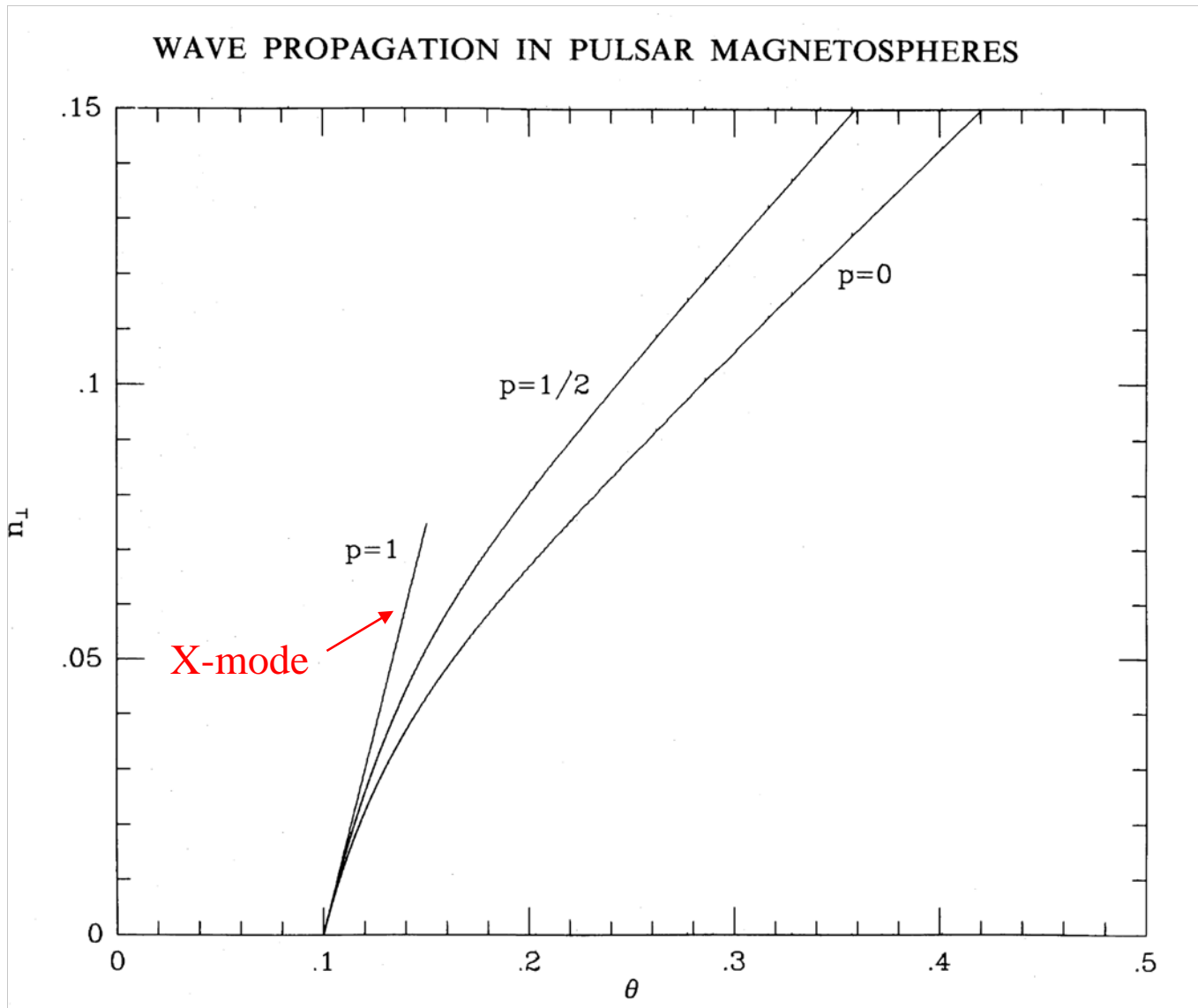


Аномальный момент, действующий на вращающийся шар в вакууме, и другие простые задачи

Бескин В.С., Желтоухов А.А.
ФИАН, МФТИ

Refraction

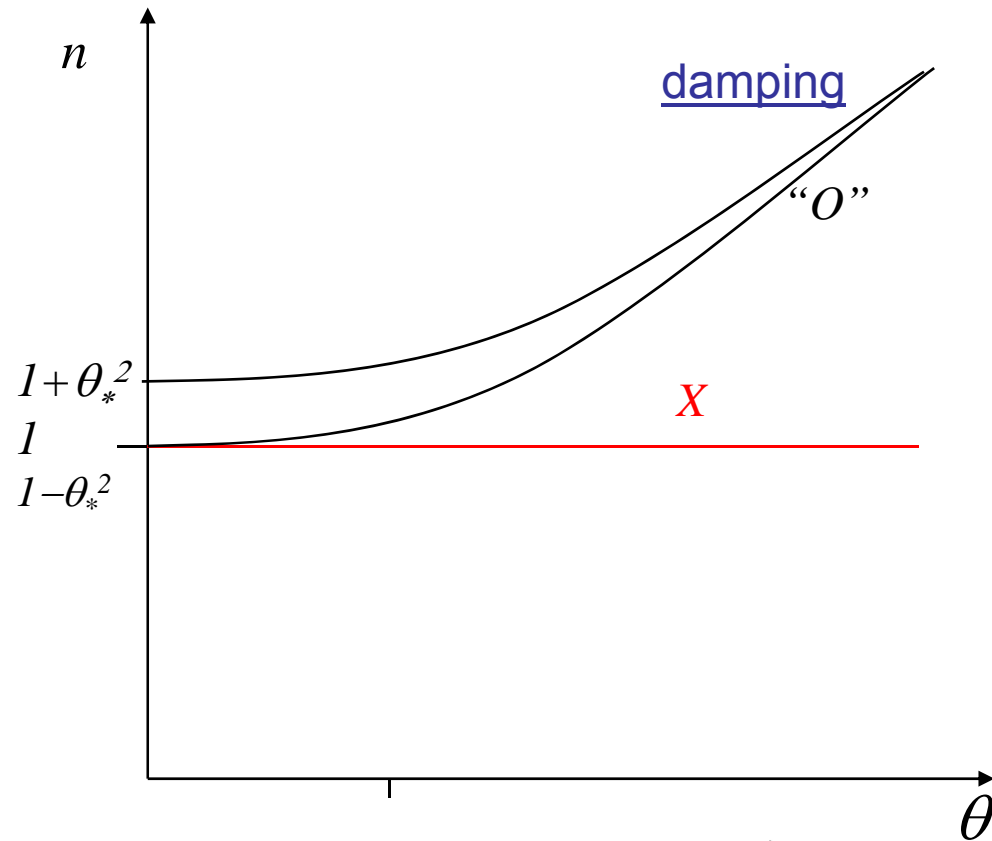
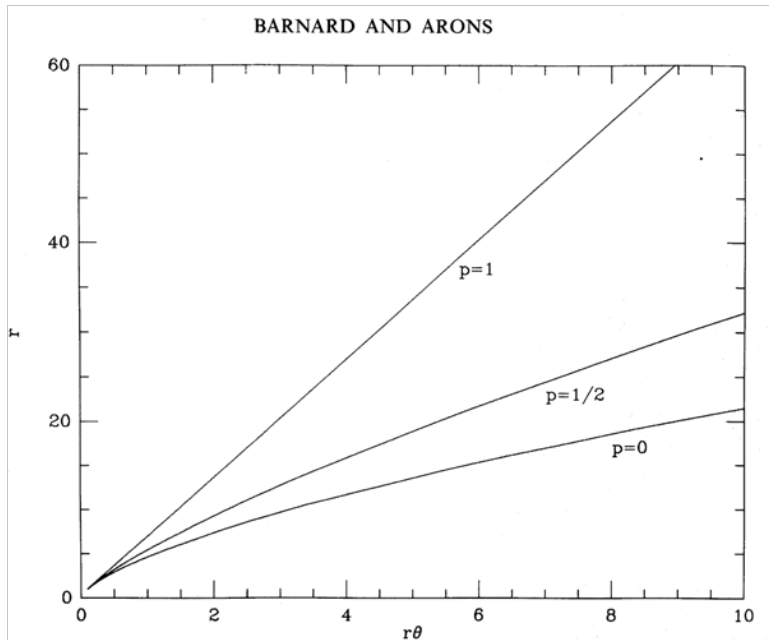
J.Arons, J.Barnard. ApJ, **302**, 120 (1986)



Refraction

J.Barnard, J.Arons, ApJ, **302**, 138 (1986)

$$\varepsilon_{ij} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 - \left\langle \frac{\omega_p^2}{\omega^2 \gamma^3} \right\rangle \end{pmatrix}$$



$$\theta_* = \left\langle \left(\frac{\omega_p^2}{\omega^2 \gamma^3} \right)^{1/4} \right\rangle$$

Refraction

VB, A.V.Gurevich & Ya.N.Istomin, Ap&SS, **146**, 205 (1988)

$$\varepsilon_{ij} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 - \langle \frac{\omega_p^2}{\omega^2 \gamma^3} \rangle \end{pmatrix}$$

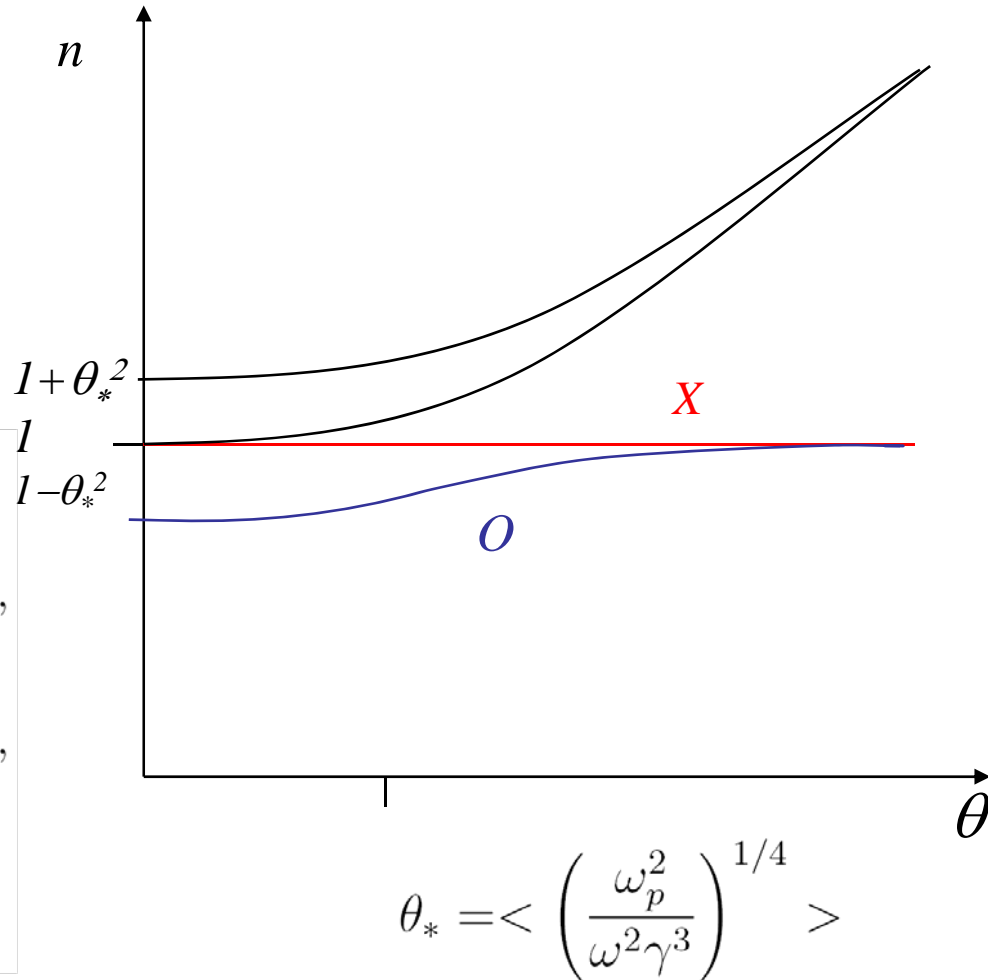
If $A_p = \frac{\omega_p^2}{\omega^2} \langle \gamma \rangle \gg 1$

$$n_1 = 1,$$

$$n_2 = 1 + \frac{\theta^2}{4} - \left(\frac{\omega_p^2}{\omega^2} \langle \frac{1}{\gamma^3} \rangle + \frac{\theta^4}{16} \right)^{1/2},$$

$$n_3 = 1 + \frac{\theta^2}{4} + \left(\frac{\omega_p^2}{\omega^2} \langle \frac{1}{\gamma^3} \rangle + \frac{\theta^4}{16} \right)^{1/2},$$

$$n_4 = \frac{1}{\cos \theta}.$$



Propagation

$$\frac{dr_{\perp}}{dl} = \frac{\partial}{\partial k_{\perp}} \left(\frac{k}{n_j} \right),$$

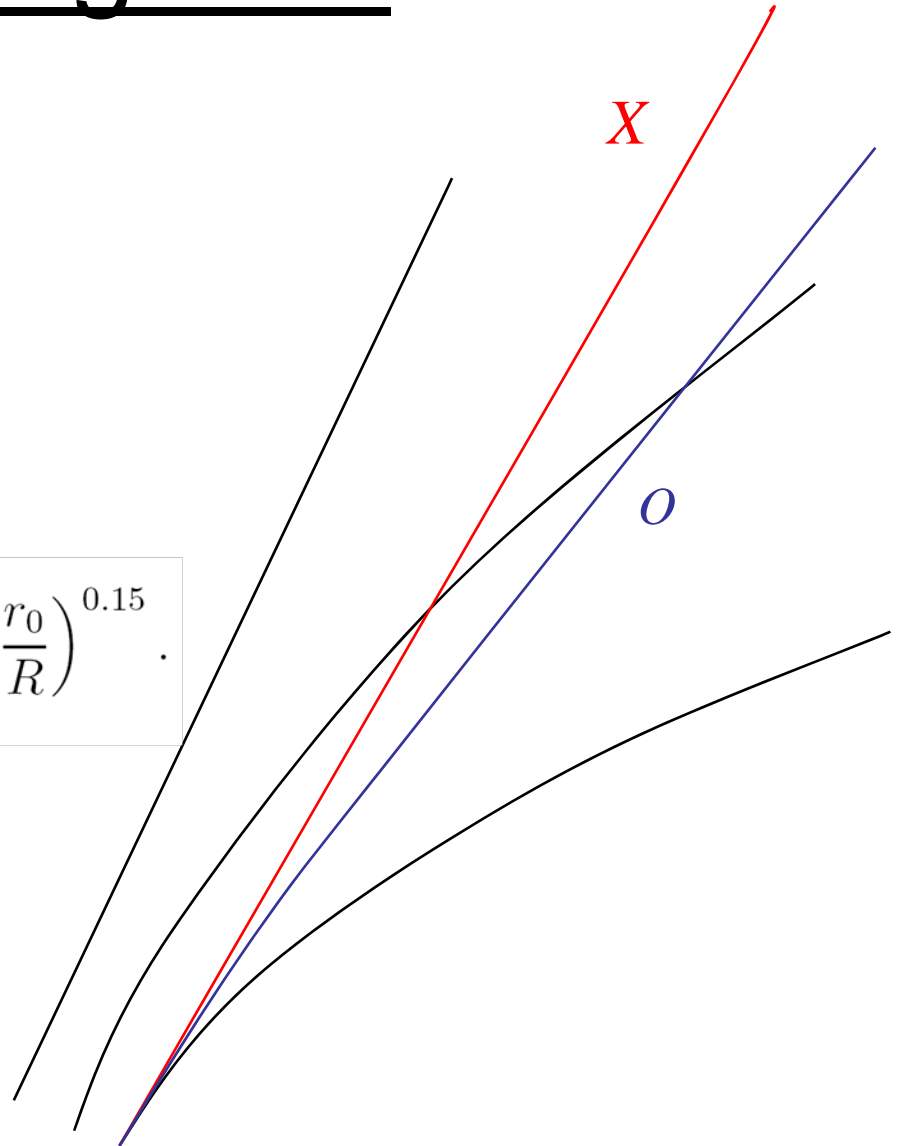
$$\frac{dk_{\perp}}{dl} = - \frac{\partial}{\partial r_{\perp}} \left(\frac{k}{n_j} \right)$$

$$W \approx \left(\frac{\Omega R}{c} \right)^{0.36} \left(\frac{\omega_{p0}}{\omega} \right)^{0.14} \langle \gamma^{-3} \rangle^{0.07} \left(\frac{r_0}{R} \right)^{0.15} .$$

On can know $r_0(v)$

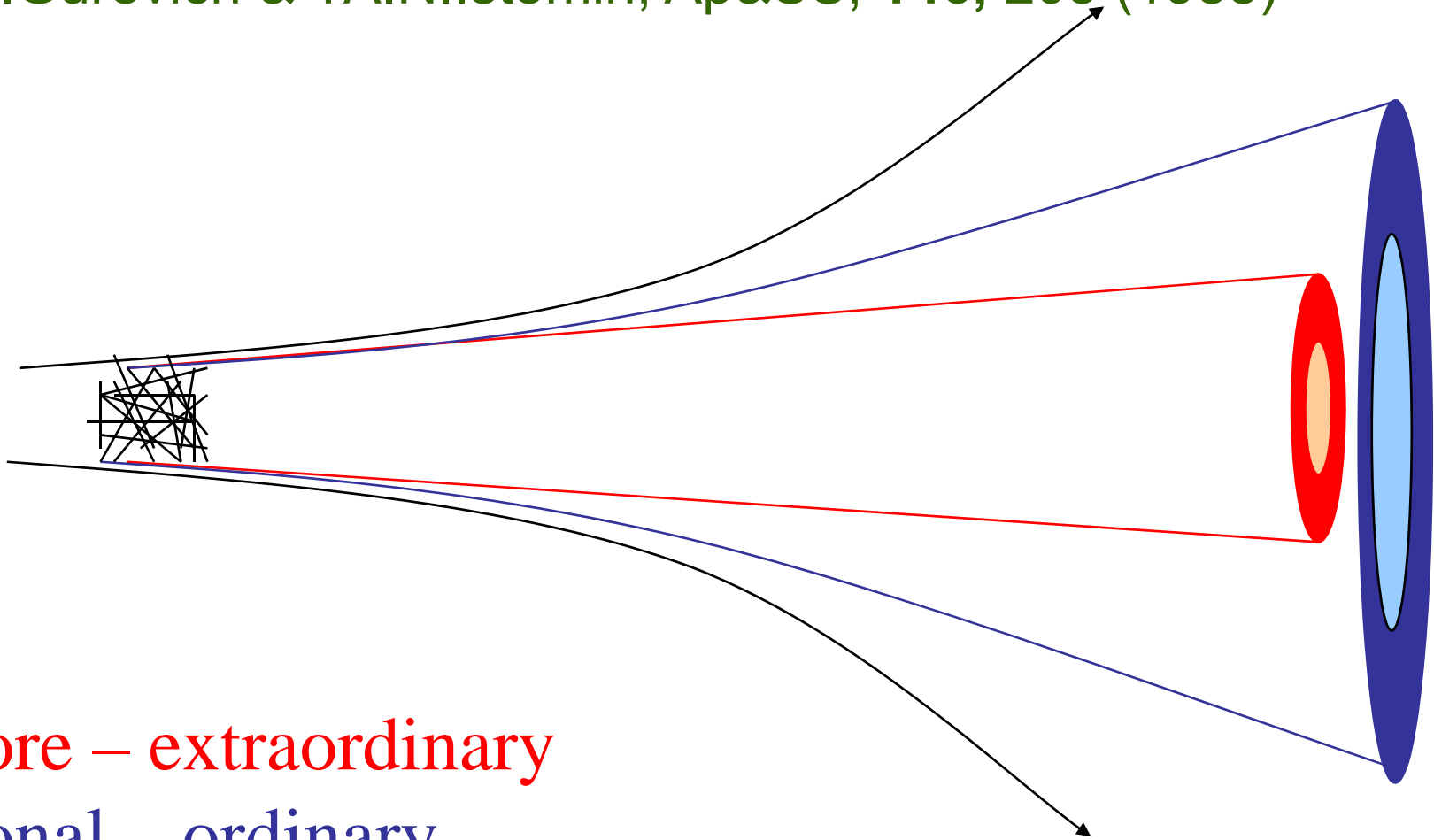
Yu. Lyubarsky, S.Petrova

D.Lai et al, VB et al



Core & Conal

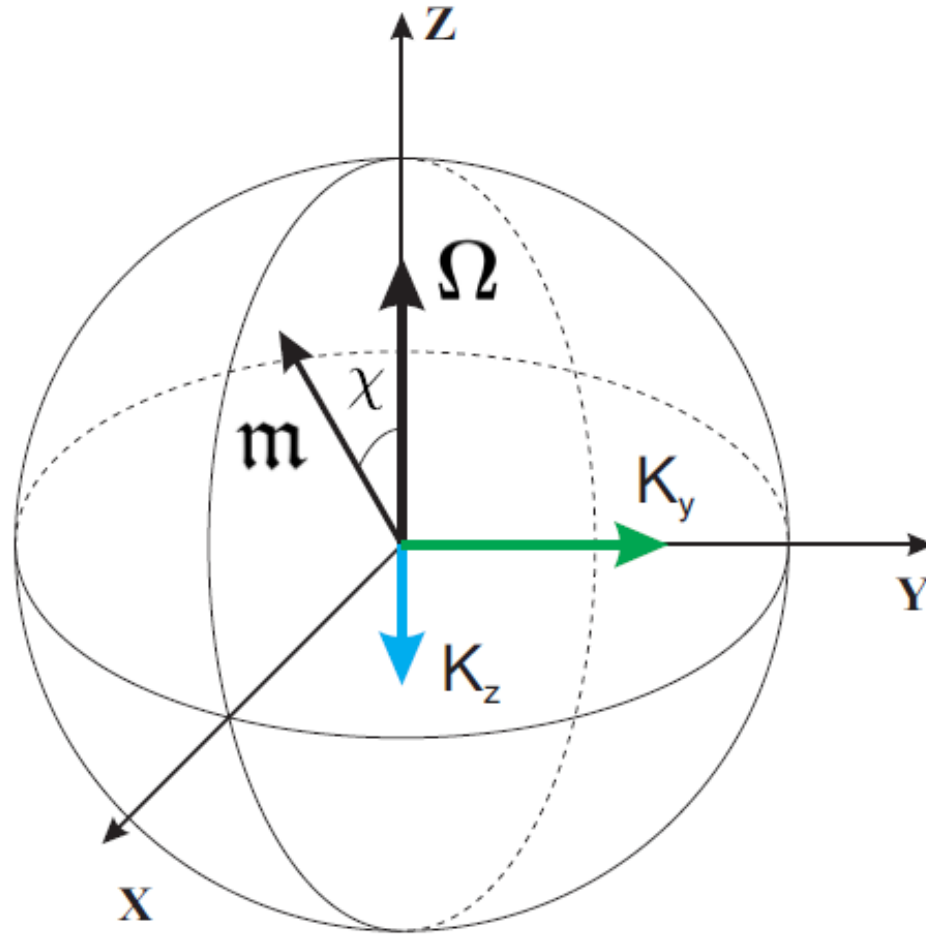
VB, A.V.Gurevich & YA.N.Istomin, Ap&SS, **146**, 205 (1988)



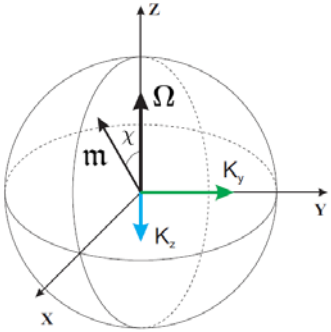
Core – extraordinary

Conal – ordinary

Anomalous torque



Anomalous torque



$$K_{z'} = \frac{2 m^2}{3 R^3} \left(\frac{\Omega R}{c} \right)^3 \sin^2 \chi$$

$$K_{x'} = \frac{2 m^2}{3 R^3} \left(\frac{\Omega R}{c} \right)^3 \sin \chi \cos \chi$$

$$K_{y'} = \xi \frac{m^2}{R^3} \left(\frac{\Omega R}{c} \right)^2 \sin \chi \cos \chi$$

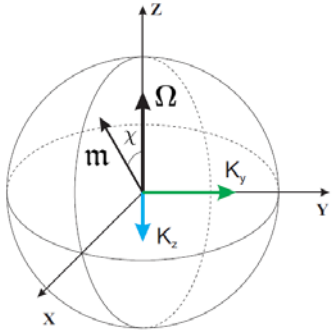
$\xi = 1$ (Davis & Goldstein 1970, Goldreich 1970)

$\xi = 1/5$ (Good&Ng 1985)

$\xi = 0$ (Michel 1991, Istomin 2005)

$\xi = 3/5$ (Melatos 2000)

Anomalous torque



Main assumptions

- Rotation in vacuum
- Freezing-in condition inside the body

$$\mathbf{E} + \beta_{\mathbf{R}} \times \mathbf{B} = 0$$

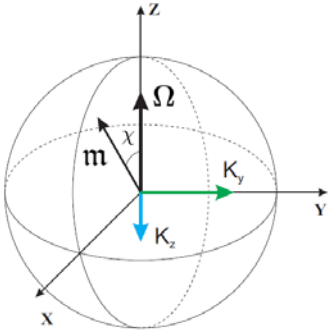
$$\beta_{\mathbf{R}} = \boldsymbol{\Omega} \times \mathbf{r} / c$$

- Corotation currents $\mathbf{j} = c\rho_e\beta_{\mathbf{R}}$ only
(hence, zero volume force)

$$d\mathbf{F} = \rho_e\mathbf{E} dV + [\mathbf{j} \times \mathbf{B}] / c dV$$

- Spherical body

Anomalous torque



Basic equation

$$\frac{d \mathbf{L}_{\text{field}}}{dt} + \mathbf{K}^M + \int [\mathbf{r} \times \mathbf{F}] dV = 0$$

Electromagnetic angular momentum

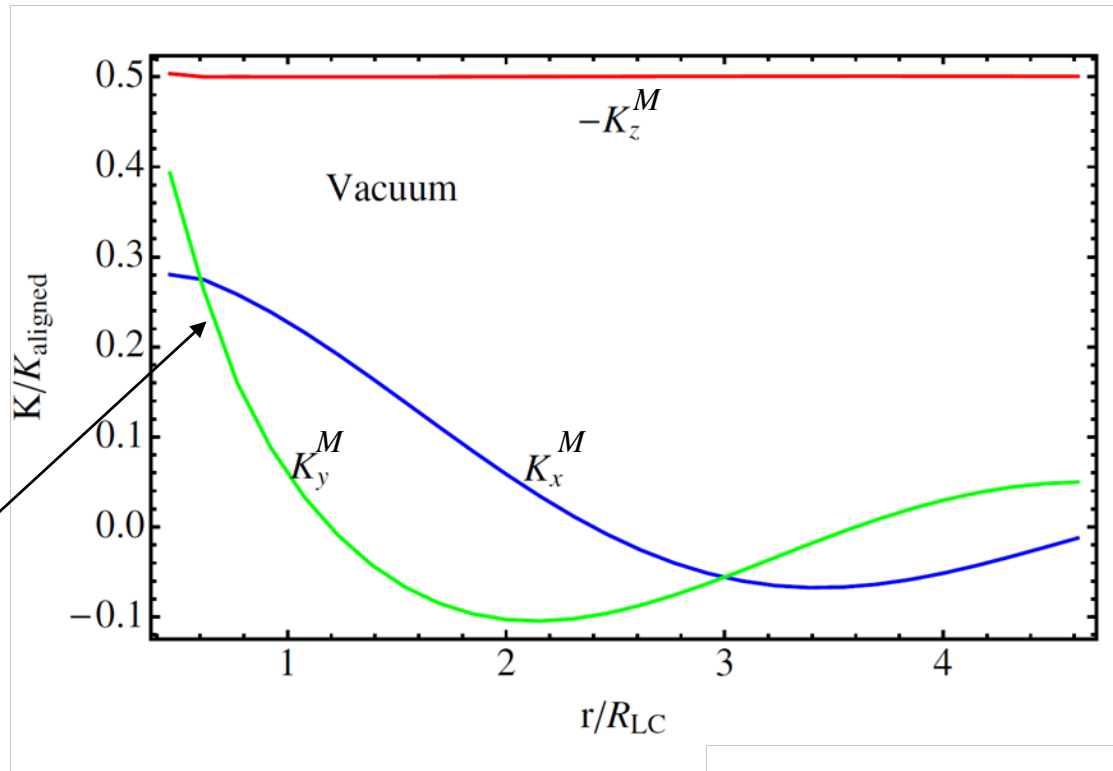
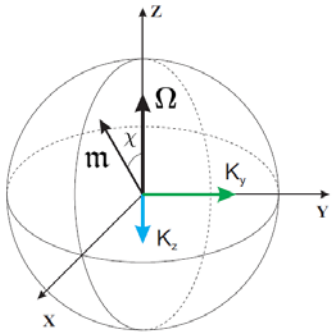
$$\mathbf{L}_{\text{field}} = \int \frac{[\mathbf{r} \times [\mathbf{E} \times \mathbf{B}]]}{4\pi c} dV$$

Flux $K_i^M = - \int \varepsilon_{ijk} r_j T_{kl} dS_l$ (spherical surface)

$$\mathbf{K}_{y'}^M = \frac{R^3}{4\pi} \int \left([\mathbf{n} \times \mathbf{B}]_{y'} (\mathbf{B}\mathbf{n}) + [\mathbf{n} \times \mathbf{E}]_{y'} (\mathbf{E}\mathbf{n}) \right) d\omega$$

Anomalous torque

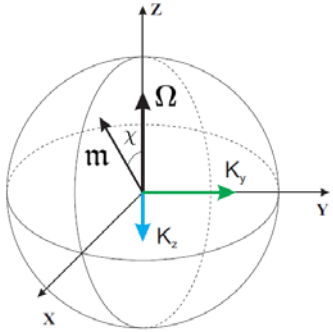
A.Philippov, A.Tchekhovskoy, J.Li
<http://arxiv.org/abs/1311.1513>



$1/r$

$$\frac{d\mathbf{L}_{\text{field}}}{dt} + \mathbf{K}^M + \int [\mathbf{r} \times \mathbf{F}] dV = 0$$

Anomalous torque



Torque, acting on the body

$$\frac{d \mathbf{L}_{\text{mat}}}{dt} = \int [\mathbf{r} \times \mathbf{F}] dV$$

$$d\mathbf{F} = \sigma_e \mathbf{E} dS + [\mathbf{I}_S \times \mathbf{B}] / c dS$$

$$\mathbf{K} = \int \mathbf{r} \times d\mathbf{F} = \frac{R^3}{4\pi} \int \left([\mathbf{n} \times \{\mathbf{B}\}] (\mathbf{B}\mathbf{n}) + [\mathbf{n} \times \mathbf{E}] (\{\mathbf{E}\} \mathbf{n}) \right) d\omega$$

- for spherical body
- for highly conducting interior
- for corotation currents only

Anomalous torque

Calculations

- small parameter
- quasi-stationary

$$\varepsilon = \frac{\Omega R}{c}$$

$$\varphi - \Omega t$$

$$\nabla \times (\mathbf{E} + \beta_R \times \mathbf{B}) = 0,$$

$$\nabla \times (\mathbf{B} - \beta_R \times \mathbf{E}) = \frac{4\pi}{c} \mathbf{j} - 4\pi \rho_e \beta_R$$

$$\beta_R = \boldsymbol{\Omega} \times \mathbf{r} / c$$

$$\mathbf{j} = c \rho_e \beta_R$$

$$\mathbf{E} + \beta_R \times \mathbf{B} = -\nabla \psi,$$

$$\mathbf{B} - \beta_R \times \mathbf{E} = \nabla h$$

$$\longrightarrow \psi^{(1)} \longrightarrow E^{(1)}$$

$$\longrightarrow h^{(2)} \longrightarrow B^{(2)}$$

Anomalous torque

$$\begin{aligned} B_r^\perp &= \frac{|\mathbf{m}|}{r^3} \sin \theta \operatorname{Re} \left(2 - 2i \frac{\Omega r}{c} \right) \exp \left(i \frac{\Omega r}{c} + i\varphi - i\Omega t \right), \\ B_\theta^\perp &= \frac{|\mathbf{m}|}{r^3} \cos \theta \operatorname{Re} \left(-1 + i \frac{\Omega r}{c} + \frac{\Omega^2 r^2}{c^2} \right) \exp \left(i \frac{\Omega r}{c} + i\varphi - i\Omega t \right), \\ B_\varphi^\perp &= \frac{|\mathbf{m}|}{r^3} \operatorname{Re} \left(-i - \frac{\Omega r}{c} + i \frac{\Omega^2 r^2}{c^2} \right) \exp \left(i \frac{\Omega r}{c} + i\varphi - i\Omega t \right), \\ E_r^\perp &= 0, \\ E_\theta^\perp &= \frac{|\mathbf{m}| \Omega}{r^2 c} \operatorname{Re} \left(-1 + i \frac{\Omega r}{c} \right) \exp \left(i \frac{\Omega r}{c} + i\varphi - i\Omega t \right), \\ E_\varphi^\perp &= \frac{|\mathbf{m}| \Omega}{r^2 c} \cos \theta \operatorname{Re} \left(-i - \frac{\Omega r}{c} \right) \exp \left(i \frac{\Omega r}{c} + i\varphi - i\Omega t \right). \end{aligned}$$

Anomalous torque

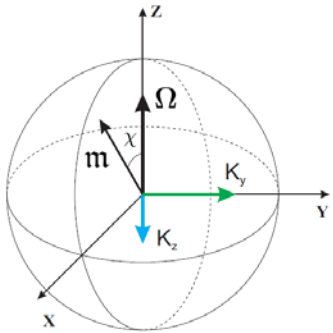
$$\mathbf{E} + \beta_{\text{R}} \times \mathbf{B} = -\nabla\psi,$$

$$\mathbf{B} - \beta_{\text{R}} \times \mathbf{E} = \nabla h$$

Calculations

- Freezing-in condition $\psi^{(\text{In})} = 0$
- Hence, no freedom in $E^{(1)}$
- Exactly corresponds to vacuum solution
- Freedom in $B^{(2)}$ (free functions $h^{(2)}$)
- Result DOES NOT depend on this freedom

Anomalous torque



Torque, acting on the body

$$\mathbf{K} = \int \mathbf{r} \times d\mathbf{F} = \frac{R^3}{4\pi} \int \left([\mathbf{n} \times \{\mathbf{B}\}] (\mathbf{B}\mathbf{n}) + [\mathbf{n} \times \mathbf{E}] (\{\mathbf{E}\} \mathbf{n}) \right) d\omega$$

$$K = K^M(R + 0) - K^M(R - 0)$$

$$\xi = \frac{3}{5}$$

- Does not depend on the internal structure and freedom in $h^{(2)}$
- Obtained by A.Melatos (2000) for Deutsch solution

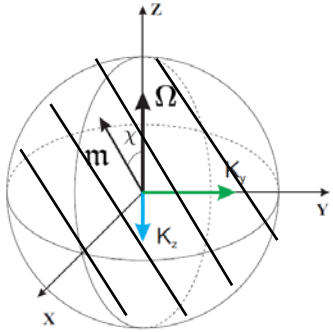
$$\mathbf{K}^M = - \frac{d\mathbf{L}_{\text{field}}}{dt}$$

$$\mathbf{L}_{\text{field}} = \int \frac{[\mathbf{r} \times [\mathbf{E} \times \mathbf{B}]]}{4\pi c} dV$$

Depends on the internal structure

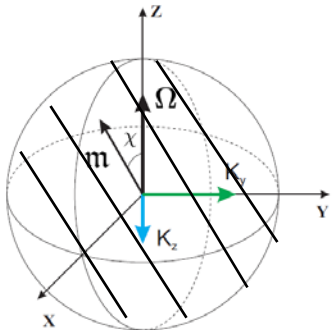
Anomalous torque

Three cases



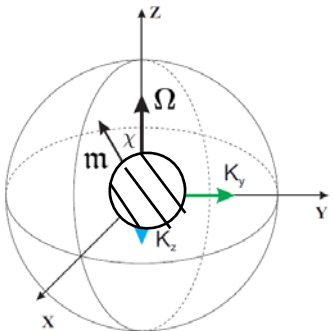
- Homogeneously magnetized sphere

$$K_{y'} = \frac{1}{3} \frac{m^2}{R^3} \left(\frac{\Omega R}{c} \right)^2 \sin \chi \cos \chi$$



- Homogeneously magnetized hollow envelope

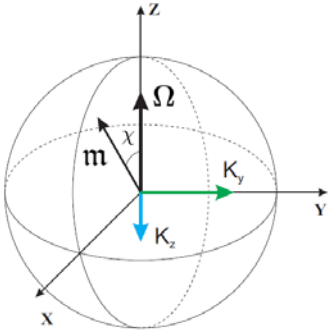
$$K_{y'} = \frac{31}{45} \frac{m^2}{R^3} \left(\frac{\Omega R}{c} \right)^2 \sin \chi \cos \chi$$



- Homogeneously magnetized core $R_{\text{in}} < R$

$$K_{y'} = \left(\frac{8}{15} - \frac{1}{5} \frac{R}{R_{\text{in}}} \right) \frac{m^2}{R^3} \left(\frac{\Omega R}{c} \right)^2 \sin \chi \cos \chi$$

Anomalous torque

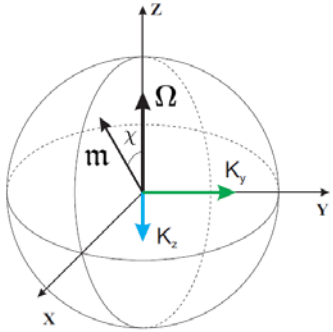


Main conclusion

Anomalous torque acting on the body depends on the internal structure of the electromagnetic field inside the body:



Anomalous torque



More conclusions

- Nonzero angular momentum L_{field} of the electromagnetic field exist only in Ω^2 order.
- In Ω^3 order it vanishes. It implies no problem in calculating the braking torque.
- Hence, energy losses can be determined via

the flux
$$K_i = - \int \varepsilon_{ijk} r_j T_{kl} dS_l$$